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Observation of a Narrow Anti-Baryon State at 2.26 GeV/c² *
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ABSTRACT

We report evidence, from a study of multihadron final states produced in the wide-band photon beam at Fermilab, for the production of a new antibaryon state which decays into $\bar{h}\pi^-\pi^-\pi^+$. The mass of the state is 2.26 \pm 0.01 GeV/c² and the decay width is < 75 MeV/c². We also report evidence of a higher mass state (~ 2.5 GeV/c²) which decays into the 2.26 GeV/c² state.

This paper is based on the study of multihadron final states produced in the wide-band photon beam at Fermilab. We report on the observation of a narrow peak near 2.26 GeV/c² in the invariant mass spectrum of the $\Lambda \pi \pi \pi \pi^+$ final state. We do not observe a significant peak near 2.26 GeV/c² in the $\Lambda \pi^+ \pi^+ \pi^-$ final state. We also observe evidence that there exists a state of higher mass near 2.5 GeV/c² which then decays into $\Lambda \pi \pi \pi^+ \pi^+$ (2.26 GeV) + π^+ .

Our results are based on studies of multihadron data taken during September - December of 1975 in the wide-band photon beam at Fermilab. The results of similar studies with data taken earlier have been previously reported. A description of the detector can be found in Refs. 1 and 2.

All pairs of oppositely charged tracks which verticize more than 27 in. downstream of the production target are potential candidates For $K_S \to \pi^+\pi^-$, $\Lambda(\bar{\Lambda}) \to p(\bar{p})\pi^-(\pi^+)$, or new long-lived objects. For events considered in this paper, we further require that there be a vertex within the production target with more than two charged tracks, and that the v^0 intersect this vertex. If the invariant mass of the v^0 is the mass of K_S (498 \pm 15 MeV/ c^2), it is not considered as a candidate for Λ or $\bar{\Lambda}$. (We lose \sim 10% of Λ due to K_S/Λ mass ambiguity.) For the remaining v^0 's, we plot the invariant mass assuming tracks are $\pi^-p(\pi^+\bar{p})$ and they are shown in Figs. la and lb. We make a mass cut of 1116 \pm 2 MeV/ c^2 . The background to the $\Lambda(\bar{\Lambda})$ signal is a few percent.

We have studied the invariant mass distribution of $\overline{h}nv^{\pm}$ (n = 1,2,3,4) in all charge states. In this paper, we will discuss n = 3 and n = 4 states only. In order to be included in our studies, events must pass the following selection criteria: 1) The total energy of the charged tracks of an event is less than 200 GeV. 2) The total number of reconstructed tracks is less than nine. These cuts reduce neutron-induced events.

In Figs. 2a and 2b, the invariant mass distributions of all combinations of $\Lambda \pi \pi \pi^{\dagger} \pi^{\dagger}$ and $\Lambda \pi \pi^{\dagger} \pi^{\dagger} \pi^{\dagger}$ are shown. We observe a clear peak near 2.26 GeV/c² in $\Lambda \pi \pi^{\dagger} \pi^{\dagger}$, but not in $\Lambda \pi^{\dagger} \pi^{\dagger} \pi^{\dagger}$ final states. The statistical significance of the peak is ~ 7 standard deviations. We estimate the mass of the peak to be 2.26 \pm 0.01 GeV/c² and the measured width to be 40 \pm 20 MeV/c² (FWHM). The measured width is consistent with our experimental mass resolution (~ 30 MeV/c²) for a zero-width state.

We have also studied the invariant mass distributions of the states $\Lambda \pi^+ \pi^-$ and $\Lambda \pi^- \pi^- \pi^+$. At present, we neither observe nor rule out a similar signal near 2.26 GeV in either state. The number of $\Lambda 3\pi^{\pm}$ combinations in the vicinity of the mass of 2.26 GeV is about three times larger than that of $\Lambda 3\pi^{\pm}$, a result, no doubt, of the excess of baryons over antibaryons in the neutral beam and the production target.

The first question of interest is whether this is a well-known hadron or a new member of a well-known hadron family. There exists in the literature a particle, $\Sigma(2.25~{\rm GeV/c^2}), \ {\rm with\ a\ full\ width\ of\ 100-230\ MeV/c.}^{2,3}$ The upper limit of our measured width is substantially smaller than the quoted width. We also believe that if this were an ordinary hadron, we should have observed a peak in $\Lambda\pi^+\pi^+\pi^-$.

One of the attractive features of the $\overline{\Lambda}$ final state is that, by studying the longitudinal polarization of $\overline{\Lambda}$, it is possible to determine whether or not parity is violated. We studied, in the $\overline{\Lambda}3\pi$ rest system, the forward/backward asymmetry in the angular distribution of the \overline{p} with respect to the $\overline{\Lambda}$ direction. Although we do not observe a statistically significant asymmetry, we do find that the \overline{p} favors the forward direction. The effect is $\sim 2~\sigma$. We also look at the energy distribution of the π^+ and the π^- in the $\overline{\Lambda}\pi^-\pi^-$ state. On the average, the energy of the π^+ is higher than that of the π^-

We also observe evidence of a cascade $\bar{\Lambda}(4\pi)^{\circ}$ (~ 2.5 GeV/c²) $\rightarrow \bar{\Lambda}(3\pi)^{-1}$ (2.26) $+\pi^{+}$, which we illustrate in Fig. 3. Invariant masses for all combinations of $\bar{\Lambda}(4\pi)^{\circ}$ are shown in Fig. 3a. For each of these entries, we then plot the masses of both $\bar{\Lambda}(3\pi)^{-}$ combinations in Fig. 3b. In Fig. 3c, we plot the mass difference between the $\bar{\Lambda}(4\pi)^{\circ}$ and

 $\overline{\Lambda}$ (3 π)⁻, separately for $\overline{\Lambda}$ (3 π)⁻ mass within the 2.26 peak (2.25-2.275) and immediately outside (dashed). A clear peak for mass difference \sim 200 MeV/c² is visible only in association with the $\overline{\Lambda}$ (3 π)⁻ peak at 2.26 GeV/c².

In summary, we observe a clear peak near 2.26 GeV in the $\overline{\Lambda}\pi^+\pi^-\pi^+\pi^+$ final state, though not in the $\overline{\Lambda}\pi^+\pi^+\pi^-$ final state. The measured width of this state is consistent with the experimental mass resolution of a zero width state. We also observe evidence of the cascade process $\overline{\Lambda}\pi^+\pi^-\pi^+\pi^-$ ($\sim 2.5~{\rm GeV/c}^2$) $\rightarrow \overline{\Lambda}\pi^-\pi^-\pi^+$ (2.26 GeV/c²) $+\pi^+$. It is also noted that one neutrino event reported by Cazzoli et al has the same mass as the $\overline{\Lambda}\pi^-\pi^-\pi^+$ state and shows the similar cascade process.

Our results are consistent with predictions of a lowest mass charmed baryon (c_0^+/Λ_c^+) near 2.26 GeV/ c^2 and two next higher states (c_1/Σ_c , c_1^+/γ_c^+) near 2.42 and 2.48 GeV/ c^2 , which cascade to c_0^+/Λ_c^+ π . Together with the narrow mesonic states recently observed by the SLAC-LBL group, and one unusual neutrino event observed at BNL, our observation lends further evidence for the existence of a new quantum number, charm. 7,8

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- 1 B. Knapp et al, Phys. Rev. Lett. 34, 1040 (1975).
- W. Lee, Proceedings of the 1975 International Symposium, Stanford, Calif., p. 213.
- N. Samios et al, "Hadrons and SU 3", BNL Preprint.

 R. Cool et al, Phys. Rev. Dl, 1887 (1970).
- ⁴ E.G. Cazzoli et al, Phys. Rev. Lett. <u>34</u>, 1125 (1975).
- 5 A. DeRujula, H. Georgi, S.L. Glashow, Harvard Preprint.
- 6 G. Goldhaber et al, Phys. Rev. Lett. 37, 255 (1976).

 I. Peruzzi et al, to be published.
- 7 B.J. Bjorken, S.L. Glashow, Phys. Lett. <u>11</u>, 255 (1964);
 S.L. Glashow, J. Iliopoulos, L. Maiani, Phys. Rev. <u>D1</u>,
 1285 (1970); M.K. Gaillard, B.W. Lee, J.L. Rosner,
 Rev. Mod. Phys. <u>47</u>: 277 (1975); A. DeRujula, H. Georgi,
 S.L. Glashow, Phys. Rev. <u>D12</u>, 147 (1975).
- Other indications of possible charm particle production have come from neutrino induced dilepton experiments.

 See, for example, A. Benvenuti et al, Phys. Rev. Lett. 34,
 419 (1975); J. Bleitschaff et al, Phys. Lett. 60B, 207 (1976);

 J. von Krogh et al, Phys. Rev. Lett. 36, 710 (1976);

 B.C. Barish et al, Phys. Rev. Lett. 36, 939 (1976);

 W. Lee et al and C. Baltay et al, in Proceedings of the Aachen Neutrino Conference (1976), to be published.

- Fig. 1 Invariant mass distributions for V^{O} , assuming: (a) $p\pi^{-}$; (b) $\bar{p}\pi^{+}$.
- Fig. 2 Invariant mass distributions for $\bar{\Lambda}3\pi$ combinations, separated by total charge: (a) $\bar{\Lambda}\pi^{\dagger}\pi^{\dagger}\pi^{-}$; (b) $\bar{\Lambda}\pi^{\dagger}\pi^{\dagger}\pi^{-}$.
- Fig. 3 study of cascade $\overline{\Lambda}(4\pi)^{\circ} \rightarrow \overline{\Lambda}(3\pi)^{-} + \pi^{+}$:
 - (a) Invariant mass of all $\bar{\Lambda}(4\pi)^{\circ}$ combinations;
 - (b) Invariant mass of all $\overline{\Lambda}(3\pi)^-$ combinations from $\overline{\Lambda}(4\pi)^{\circ}$;
 - (c) Mass difference between $\bar{\Lambda}(4\pi)^{\circ}$ and $\bar{\Lambda}(3\pi)^{-}$, separately for $\bar{\Lambda}(3\pi)^{-}$ within 2.26 peak and immediately outside the peak (dashed curve).











